

# The Development of Particle Physics

Dr. Vitaly Kudryavtsev  
E45, Tel.: 0114 2224531  
[v.kudryavtsev@sheffield.ac.uk](mailto:v.kudryavtsev@sheffield.ac.uk)



# CP violation

---

- Charge conjugation and time reversal symmetries.
- CPT theorem.
- Neutral kaon system.
- CP violation in neutral kaon system.
- Experimental evidence for CP violation in neutral kaon system.
  - Experiment
  - Data analysis
  - Results and outcomes
- Other possible examples of CP violation.
- Summary and conclusions.

# Charge conjugation

- **C operation - interchange of particle with its antiparticle.**
- C symmetry in classical physics - invariance of Maxwell's equations under change in sign of the charge, electric and magnetic fields.
- C symmetry in particle physics - the same laws for a set of particles and their antiparticles: collisions between electrons and protons are described in the same way as collisions between positrons and antiprotons. The symmetry also applies to neutral particles.
- **$C\psi = \pm \psi$ : even or odd symmetry.**
- Example: particle decay into two photons, for example  $\pi^0 \rightarrow 2\gamma$ , by the electromagnetic force. Photon is odd under C symmetry; two photon state gives a product  $(-1)^2$  and is even. So, if symmetry is exact, then 3 photon decay is forbidden. In fact it has not been observed.
- **C symmetry holds in strong and electromagnetic interactions.**

# Charge conjugation and time reversal.

- As we saw, C invariance was violated in weak interactions because parity was violated, if CP symmetry was assumed to be preserved.
- Under C operation left-handed neutrinos should transform into left-handed antineutrino, which was not found in nature. However, the combined CP operation transforms left-handed neutrino into right-handed antineutrino, which exists.
- **T operation - connect a process with that obtained by running backwards in time: reverses the directions of motion within the system.**
- **T symmetry: "initial state  $\rightarrow$  final state" can be converted to "final state  $\rightarrow$  initial state" by reversing the directions of motion of all components of the system.**

# CPT theorem

- It is possible to define product symmetries, like CP (parity and charge conjugation simultaneously to reveal a system of antiparticles in the reverse-handed coordinate system) or CPT symmetry.
- **Combined CPT symmetry is absolutely exact:** for any process, its mirror image with antiparticles and time-reversal process should look exactly as the original - CPT theorem.
- **Consequences: the same masses, lifetimes, electric charges (with opposite signs) and magnetic moments for particles and antiparticles.**
- If any one individual (or pair) of the symmetries is broken, there must be a compensating asymmetry in the remaining operation(s) to ensure exact symmetry under CPT operation.
- The CPT symmetry was checked through the possible difference in masses, lifetimes, electric charges and magnetic moments of particle and antiparticles and was found to be **exact down  $10^{-19}$**  (relative difference in masses).

# Neutral kaon system

- Two isospin doublets ( $I_3 = \pm 1/2$ ,  $S = \pm 1$ ):  $K^+(u\bar{s})$  and  $K^0(d\bar{s})$  with  $S = +1$  and  $K^-(s\bar{u})$  and  $\bar{K}^0(s\bar{d})$  with  $S = -1$ .
- $K^0$  and  $\bar{K}^0$  are produced in strong interactions and are eigenstates of strong interactions as in strong interactions the third projection of isospin and strangeness are conserved (definite assignment of  $I_3$  and  $S$ ).
- However, they are not eigenstates of the CP operation. The linear combinations of these particles can be formed to identify the eigenstates of the CP:
  - $K^0_1 = (K^0 + \bar{K}^0) / \sqrt{2}$  with even CP symmetry:  $CP K^0_1 = + K^0_1$
  - $K^0_2 = (K^0 - \bar{K}^0) / \sqrt{2}$  with odd CP symmetry:  $CP K^0_2 = - K^0_2$
- The weak interaction acts on the states  $K^0_1$  and  $K^0_2$ , so the particles, which decay, are  $K^0_1$  and  $K^0_2$  without well-defined strangeness.

# Neutral kaon system

- The proof of the identity of  $K^0_1$  and  $K^0_2$  comes from their decays. Assume that CP is preserved in weak interactions. As  $K^0_1$  is even under CP operation, it can decay only to states which are also even, for example  $2\pi$ .  $K^0_2$  is odd under CP operation and can decay only to states which are also odd, for example  $3\pi$ .
- The difference in the decay modes results in the significant difference in the lifetimes:  
$$K^0_1 \rightarrow 2\pi, \quad \tau = 8.93 \times 10^{-11} \text{ s}$$
$$K^0_2 \rightarrow 3\pi, \quad \tau = 5.17 \times 10^{-8} \text{ s}$$
- Long-lived  $K^0_2$  was first observed by Lande et al. using a 3-GeV beam from the Brookhaven Cosmotron.

# Regeneration of $K^0_1$

- Pure  $K^0_2$  beam passing through matter would regenerate a  $K^0_1$  component.
- $K^0 = (K^0_1 + K^0_2) / \sqrt{2}$ ,  $\underline{K}^0 = (K^0_1 - K^0_2) / \sqrt{2}$
- Suppose we start with  $K^0$  beam, which is a combination of  $K^0_1$  and  $K^0_2$ .
- $K^0_1$  decay quickly and only  $K^0_2$  remain.  $K^0_2$  is a combination of  $K^0$  and  $\underline{K}^0$  in equal proportion.
- $K^0$  and  $\underline{K}^0$  interact differently with matter. For example, the process  $\underline{K}^0 + p \rightarrow \pi^+ + \Lambda$  is allowed, while  $K^0 + p \rightarrow \pi^+ + \Lambda$  is forbidden.
- So, after some thickness of matter the transformation will be:  
$$K^0_2 = (K^0 - \underline{K}^0) / \sqrt{2} \rightarrow (fK^0 - \underline{f}\underline{K}^0) / \sqrt{2} = (f + \underline{f}) K^0_2 / 2 + (f - \underline{f}) K^0_1 / 2$$
- As  $f \neq \underline{f}$ , the fraction of  $K^0_1$  is not zero anymore.
- **Regeneration** of  $K^0_1$  in  $K^0_2$  beam passing through matter, predicted by Pais and Piccioni, was observed by Muller et al. at the Bevatron. This allowed to determine a tiny difference in masses of  $K^0_1$  and  $K^0_2$ :  $\Delta m/m = 7 \times 10^{-15}$ .

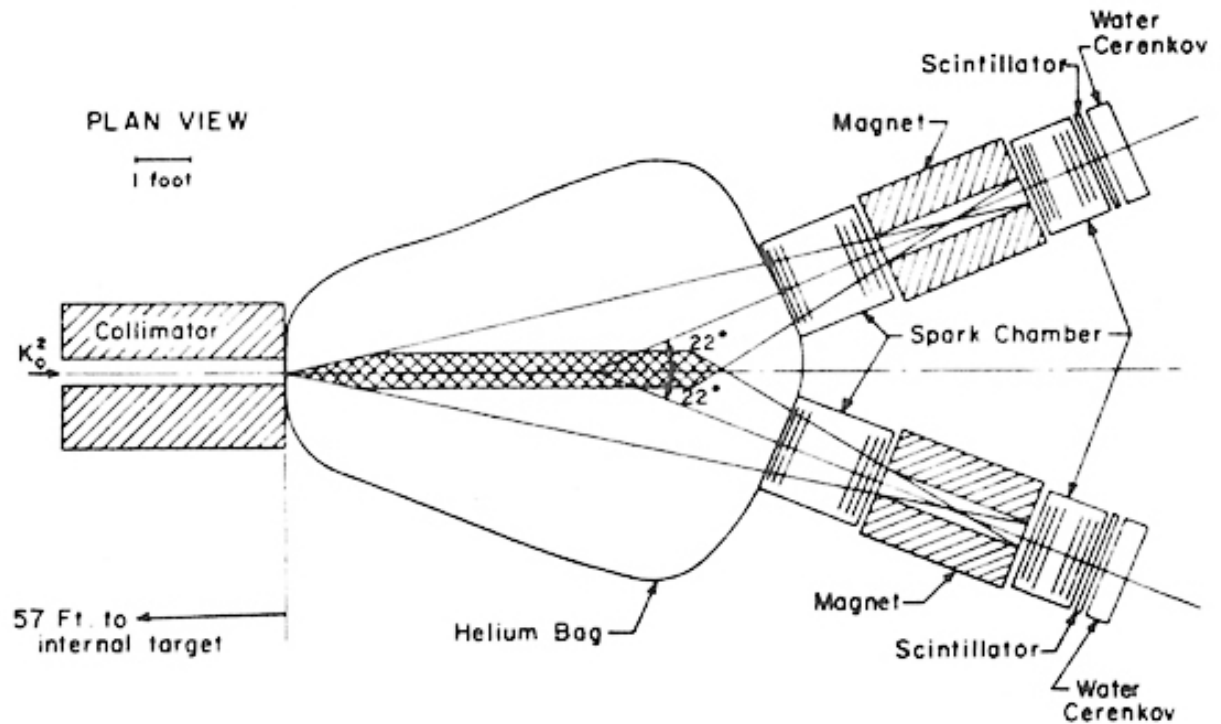


# Observation of $K^0_2 \rightarrow 2\pi$

- If  $CP$  is a good symmetry, then  $K^0_2 \rightarrow \pi^+ + \pi^-$  decay is strictly forbidden.
- Christenson et al. set up an experiment to check this. Actually, the objective of the experiment was to extend the limit on the fraction of  $K^0_2$ , which decay into two pions.
- Alternating Gradient Synchrotron (AGS) at Brookhaven.
  - 30 GeV proton beam.
  - Be target.
  - Neutral kaons were produced in  $p + Be$  collisions.
  - $K^0_2$  beam - at  $30^\circ$  relative to circulating protons.
  - Collimator at 4.5 m from the target, magnet at 6.5 m, 2nd collimator at 18 m.
- $K^0_1$  (short-lived) decayed before reaching the 2nd collimator.

# Observation of $K^0_2 \rightarrow 2\pi$

- Two spectrometers:
  - Spark chamber,
  - Magnet,
  - Scintillator,
  - Water Cherenkov counter.
- Spark chambers were triggered on a coincidence between water Cherenkov ( $v > 0.75 c$  - pions) and scintillation counters. This removed most slow particles produced in collisions of neutrons.

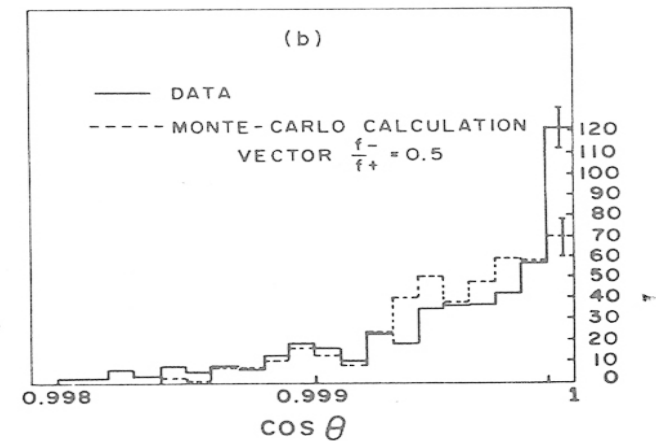
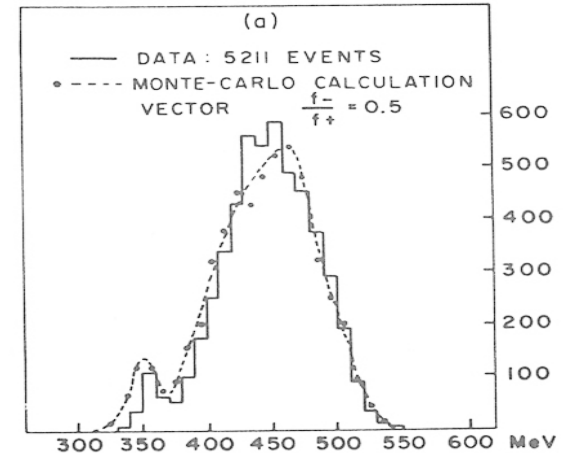


# Data analysis

- Signal events:
  - pair of particles (one particle in each spectrometer) with opposite charges;
  - invariant mass  $M_{12}^2 = (p_1 + p_2)^2 = (\mathbf{p}_1 + \mathbf{p}_2)^2 - (E_1 + E_2)^2$  should correspond to  $K^0_2$  mass (498 MeV);
  - momenta should lie in the direction of  $K^0_2$  beam.
- Background events:
  - $K^0_2 \rightarrow 3\pi$ , invariant mass of two particle system assuming each charged particle has the mass of the charged pion: 280-363 MeV,
  - $K^0_2 \rightarrow \pi \mu \nu$  (280-516 MeV),
  - $K^0_2 \rightarrow \pi e \nu$  (280-536 MeV).
- Analysis program:
  - momenta of the two detected charged particles,
  - their invariant mass assuming detected particles are charged pions,
  - the angle  $\theta$  between the vector sum of the two momenta and the direction of the  $K^0_2$  beam (zero for two-pion decay, non-zero for three-particle decay with only two particles detected).

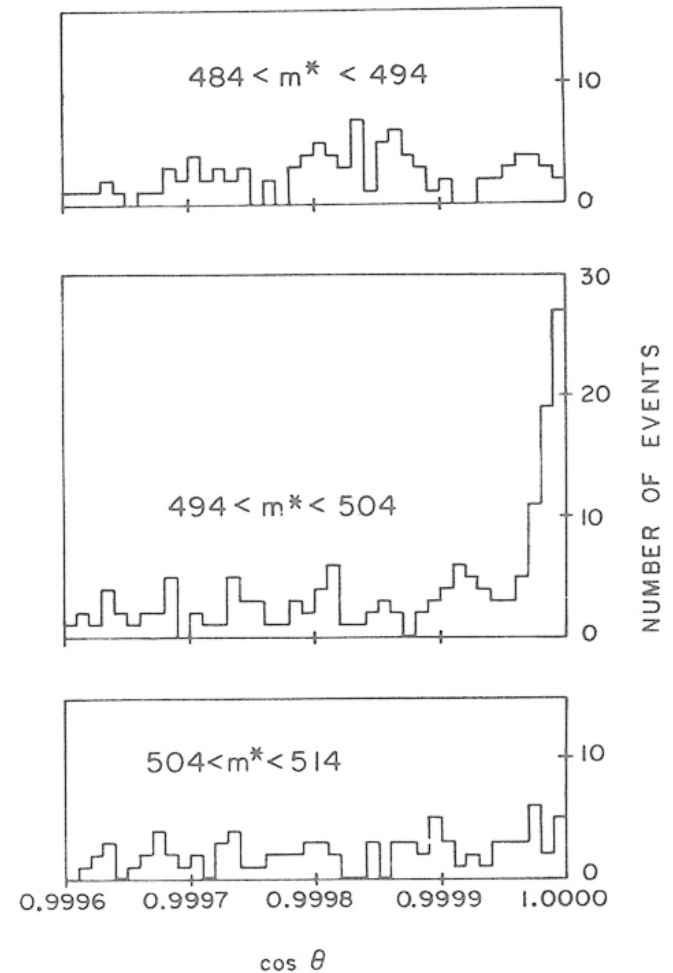
# Results

- Calibration: decays of  $K^0_1$  (into two pions) produced by regeneration in tungsten.
- Distribution in the invariant mass for all  $K^0_2$  decays (top figure).
- Angular distribution of events in the range of invariant masses from 490 to 510 MeV (expected 498 MeV): measurements and calculations for background events (bottom figure).
- Clear excess of events at small  $\theta$ .



# Results

- Angular distributions of events in three invariant mass ranges: peak at  $\theta = 0$  is seen only for masses similar to the kaon mass (498 MeV). 45 events in the forward peak after background subtraction.
- The mass of  $K^0_2$  was found from the measurements of invariant mass as  $499.1 \pm 0.8$  MeV. Not much different from the mass of  $K^0_1$  found from calibration regeneration experiment ( $498.1 \pm 0.4$  MeV).



# Results and outcomes

- Branching ratio:  $R = (K^0_2 \rightarrow \pi^+ + \pi^-) / (K^0_2 \rightarrow \text{all charged modes}) = (2.0 \pm 0.4) \times 10^{-3}$ .
- $K^0_1$  and  $K^0_2$  (in our previous definition) are not exact eigenstates of CP symmetry and are not quite particles seen by the weak interactions. Instead, the true eigenstates of CP have a small admixture of another particle and a new scheme has been adopted:

$$K^0_L = \frac{1}{\sqrt{1 + \varepsilon^2}} (K^0_2 + \varepsilon K^0_1)$$

$$K^0_S = \frac{1}{\sqrt{1 + \varepsilon^2}} (K^0_1 - \varepsilon K^0_2)$$

$\varepsilon$  is the parameter quantifying CP violation (admixture of the 2nd particle or ‘wrong’ CP eigenstate). It is measured as  $2.3 \times 10^{-3}$ .

- Similar branching ratio was found for  $K^0_L \rightarrow \pi^0 + \pi^0$  decay ( $9.1 \times 10^{-4}$ ).

# Results and outcomes

- At present the degree of CP violation is usually quoted as the amplitude ratio of processes:

$$|\eta_{+-}| \equiv \frac{\text{ampl}(K_L^0 \rightarrow \pi^+ \pi^-)}{\text{ampl}(K_S^0 \rightarrow \pi^+ \pi^-)} = (2.29 \pm 0.02) \times 10^{-3}$$

$$|\eta_{00}| \equiv \frac{\text{ampl}(K_L^0 \rightarrow \pi^0 \pi^0)}{\text{ampl}(K_S^0 \rightarrow \pi^0 \pi^0)} = (2.28 \pm 0.02) \times 10^{-3}$$

- CP violation was also demonstrated in the leptonic modes of  $K_L^0$  decay  
 $K_L^0 \rightarrow \pi^- + e^+ + \nu_e$ ,  $K_L^0 \rightarrow \pi^+ + e^- + \bar{\nu}_e$  with the asymmetry:

$$\Delta = \frac{\text{rate}(K_L^0 \rightarrow \pi^- e^+ \nu_e) - \text{rate}(K_L^0 \rightarrow \pi^+ e^- \bar{\nu}_e)}{\text{rate}(K_L^0 \rightarrow \pi^- e^+ \nu_e) + \text{rate}(K_L^0 \rightarrow \pi^+ e^- \bar{\nu}_e)} = (3.27 \pm 0.012) \times 10^{-3}$$

- CP violation provides a definition of matter and antimatter: positron (antimatter) is defined as the lepton which is produced more often in  $K_L^0$  leptonic decay. Can also separate left from right using neutrinos in  $K_L^0$  decay.

# Direct and indirect CP violation

- **Indirect CP violation:** the CP-forbidden  $K^0_1$  component in  $K^0_L$  decays via the CP-allowed process with a probability  $|\varepsilon|^2/(1+|\varepsilon|^2)\approx|\varepsilon|^2$  of finding a  $K^0_1$  component in  $K^0_L$  - process described as an admixture of the ‘wrong’ CP eigenstate.
- **Direct CP violation:** the CP-allowed  $K^0_2$  component in  $K^0_L$  decays via the CP-forbidden process:  $K^0_2 \rightarrow 2\pi$ .
- A detailed analysis of data showed that the former process dominates.
- Several experiments at CERN and Fermilab found a contribution of direct CP violation as  $\varepsilon'/\varepsilon=(1.8\pm 0.4)\times 10^{-3}$ . If this result is correct, then direct CP violation is found.
- ‘Superweak’ force, proposed by Wolfenstein, responsible for interactions with  $\Delta S=2$ . This predicts  $\varepsilon'=0$ .



# Other possible examples of CP violation

- Other possible neutral particle systems with CP violation effects:
  - $D^0$ - $\underline{D}^0$  - calculated to be very small and possibly unobservable.
  - $B^0$ - $\underline{B}^0$  - direct CP violation should dominate, ideal for its search, but many decay modes, low branching ratio to decay modes common for both  $B^0$  and  $\underline{B}^0$ , about  $10^{10}$   $B$ -meson decays are needed to measure CP violation.
  - In fact CP-violation in  $B$ -meson decays has been recently observed in several experiments.
- Cosmological CP violation is necessary to generate a cosmological matter-antimatter asymmetry:
  - To distinguish unambiguously matter from antimatter on a cosmic scale.
  - It provides difference in rate of decays of supermassive bosons  $X$  and  $Y$  into baryons and antibaryons (this asymmetry requires also non-conservation of baryon number).

# Summary and conclusions

- Kaon system provided the first evidence for CP violation in nature.
- CPT is an exact symmetry.
- Indirect CP violation dominates with a degree of  $2.3 \times 10^{-3}$ , but direct CP violation is also present.
- CP violation provides a definition of matter and antimatter.
- CP violation is required to generate cosmological matter-antimatter asymmetry.
- Big efforts are applied to find CP violation (direct violation in particular) in other systems ( $B^0 - \bar{B}^0$ ). The results show the presence of CP violation in agreement with Standard Model predictions.

# References

---

- Recommended textbooks.
- J. H. Christenson et al. “Evidence for the  $2\pi$  decay of the  $K^0_2$  meson”, *Phys. Rev. Lett.* **13** (1964) 138.
- S. Bennett et al. “Measurement of the charge asymmetry in the decay  $K^0_L \rightarrow \pi^\pm + e^\pm + \nu$ ”, *Phys. Rev. Lett.* **19** (1967) 993.